

# RMS MODE SIZE OF THE XFEL RADIATION\*

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## Abstract

Formula of the rms transverse mode size of the XFEL radiation is derived. Derivation uses information on the XFEL transverse coherence. Derived formula agrees well with the simulation result.

## DERIVATION

In this paper, the rms transverse mode size of the X-ray free electron laser (XFEL) radiation based on self amplified spontaneous emission (SASE) [1, 2, 3] is derived using the information on the transverse coherence of the radiation [4]. In general, the rms mode size of a coherent radiation is given by

$$\sigma_c = \sqrt{\lambda/(4\pi)Z_R}, \quad (1)$$

where  $\lambda$  is the radiation wavelength and  $Z_R$  is the Rayleigh length. Obviously, the XFEL mode size ( $\sigma_r$ ) must have the form of  $\sigma_c$  with proper  $Z_R$ .  $Z_R = 2L_G$ , where  $L_G = \lambda_u/(4\pi\sqrt{3}\rho)$  is the one dimensional gain length, has been generally used to give rough estimation for  $\sigma_r$  [5], which we denote as

$$\sigma_d = \sqrt{\frac{\lambda}{4\pi} \frac{2\lambda_u}{4\pi\sqrt{3}\rho}}. \quad (2)$$

Although  $\sigma_d$  is often used as a substitute for  $\sigma_r$ , it will be shown below that it is not. The derivation here uses the information of the transverse coherence of the XFEL radiation. We will consider only the cases in which  $\sigma_r$  is comparable to the rms electron beam size  $\sigma_x$ , not including the two extreme cases,  $\sigma_x \gg \sigma_r$ ,  $\sigma_r \gg \sigma_x$ , which were discussed in literatures [5].

Saldin et al. [6] computed the XFEL degree of transverse coherence as a function of

$$z = \frac{2\pi\epsilon}{\lambda} = \frac{2\pi\sigma_x^2}{\lambda\beta_0} \quad (3)$$

where  $\beta_0$  is the  $\beta$ -function that is optimized to give shortest saturation length. The result is shown in Fig. 1 and  $\beta_0$  is given by [6]

$$\beta_0 = \frac{11.2}{\sqrt{2}} \left( \frac{I_A}{I} \right)^{1/2} \frac{(\gamma\epsilon)^{3/2} \lambda_u^{1/2}}{\lambda K [JJ]}, \quad (4)$$

where  $[JJ]$  denotes the following combination of Bessel functions,

$$[JJ] = J_0 \left( \frac{K^2}{4 + 2K^2} \right) - J_1 \left( \frac{K^2}{4 + 2K^2} \right). \quad (5)$$

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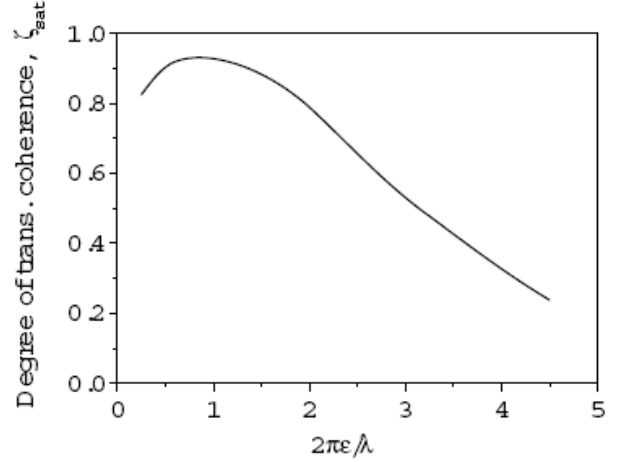


Figure 1: Graph of the degree of the transverse coherence versus  $z = 2\pi\epsilon/\lambda$ . (Ref. 6)

To make use of Fig. 1 to derive  $\sigma_r$ , we fix  $\lambda$  to a hard X-ray value. Then,  $z$  is proportional to  $\sigma_x^2$ . As  $z$  grows  $\sigma_x$  grows while  $\sigma_r$  would not grow much since  $\lambda$  is fixed. Hence, when  $z$  is large (far right side of Fig. 1),  $\sigma_x$  must be much larger than  $\sigma_r$  and the degree of transverse coherence is low. When  $z$  is very small (far left side of Fig. 1),  $\sigma_x$  must be much smaller than  $\sigma_r$  and the degree of transverse coherence is low again. The peaked part of the graph must correspond to that  $\sigma_x$  is nearly same as  $\sigma_r$ . Note, from Fig. 1, that the maximal transverse coherence is achieved at

$$\frac{2\pi\sigma_x^2}{\lambda\beta_0} = 0.8. \quad (6)$$

It is quite reasonable to assume that this maximal transverse coherence is achieved when  $\sigma_r = \sigma_x$ . Therefore, we see that

$$\sigma_r^2 = \sigma_x^2 = 0.8 \times \frac{\lambda\beta_0}{2\pi}. \quad (7)$$

To compare this expression with  $\sigma_d$ , note that the FEL parameter  $\rho$  is given by

$$\begin{aligned} \rho &= \left[ \frac{1}{8\pi} \frac{I}{I_A} \left( \frac{K[JJ]}{1 + K^2/2} \right)^2 \frac{\gamma\lambda^2}{2\pi\sigma_x^2} \right]^{1/3} \\ &= \left[ \frac{1}{8\pi} \left( \frac{I}{I_A} \frac{K^2[JJ]^2}{\gamma^3} \right) \frac{\lambda_u^2}{8\pi\sigma_x^2} \right]^{1/3}, \end{aligned} \quad (8)$$

where we used the resonance condition

$$\lambda = \frac{\lambda_u}{2\gamma^2} \left( 1 + \frac{K^2}{2} \right). \quad (9)$$

Using Eq. (4) and (6), we find that

$$\frac{I}{I_A} \frac{K^2 [JJ]^2}{\gamma^3} = \frac{11.2^2 \lambda_u \epsilon^3}{2\lambda^2 \beta_o^2} = \frac{62.7 \lambda_u \sigma_x^2}{\beta_o^3} \left(\frac{0.8}{2\pi}\right)^2. \quad (10)$$

Then  $\rho$  is simplified to

$$\rho = \left[ \frac{1}{(4\pi)^3} \frac{40.1}{4\pi} \left(\frac{\lambda_u}{\beta_o}\right)^3 \right]^{1/3} = \frac{1}{4\pi} \left(\frac{40.1}{4\pi}\right)^{1/3} \frac{\lambda_u}{\beta_o}, \quad (11)$$

and  $\beta_o$  is simply given by

$$\beta_o = \frac{1}{4\pi} \left(\frac{40.1}{4\pi}\right)^{1/3} \frac{\lambda_u}{\rho}. \quad (12)$$

Using this result in Eq. (7), it is straightforward to see that

$$\sigma_r^2 = 2.0 \times \frac{\lambda}{4\pi} \frac{2\lambda_u}{4\pi\sqrt{3}\rho} = 2.0\sigma_d^2. \quad (13)$$

Therefore, we conclude that

$$\sigma_r = 1.4\sigma_d. \quad (14)$$

Table 1: Parameters of the PAL-XFEL.

Electron energy (GeV)	3.7
Peak current (kA)	3
Normalized emittance ( $\mu\text{rad}$ )	1.0
Undulator period, $\lambda_u$ (cm)	1.5
Undulator parameter, $K$	1.49
Average $\beta$ function (m)	10
Radiation wavelength ( $\text{\AA}$ )	3
FEL parameter, $\rho$	$5.7 \times 10^{-4}$
1-D gain length, $L_G$ (m)	1.19
Saturation length, $L_{sat}$ (m)	45

## COMPARISON WITH SIMULATION

Although this was obtained with the condition  $\sigma_r = \sigma_x$ , it does not mean that Eq. (14) is valid only when  $\sigma_r = \sigma_x$ . It is important to note that the condition  $\sigma_r = \sigma_x$  was used only to derive the mode size using Fig. 1 and Eq. (14) is actually applicable to a wider range of  $\sigma_r$ . This is verified by simulation. As a demonstration, we consider PAL-XFEL that was under proposal at Pohang Accelerator Laboratory, Korea [7]. According to the PAL-XFEL parameters listed in Table 1, we find

$$\left(\frac{2\pi\epsilon}{\lambda}\right)_{PAL} = 2.89. \quad (15)$$

From Fig. 1, we see that the PAL-XFEL degree of transverse coherence is approximately 5.6, which is fairly off the maximal transverse coherence. It is easy to get

$$(\sigma_d)_{PAL} = \sqrt{\frac{\lambda}{4\pi} \frac{2\lambda_u}{4\pi\sqrt{3}\rho}} = 7.5 \mu\text{m}, \quad (16)$$

and

$$(\sigma_r)_{PAL} = 10.5 \mu\text{m}. \quad (17)$$

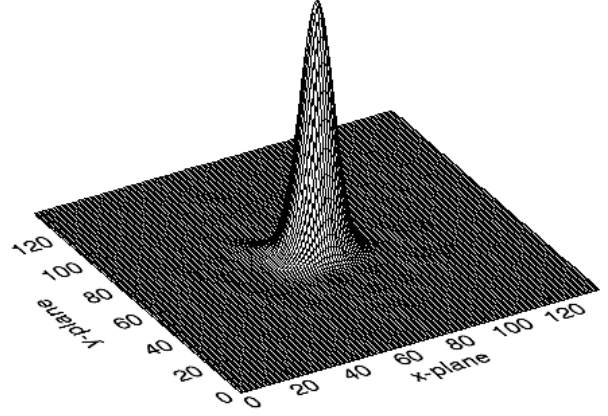


Figure 2: Simulated radiation profile of the PAL-XFEL.

The simulated radiation profile of PAL-XFEL is shown in Fig. 2 and its mode size is shown in Fig. 3. Both were simulated by GENESIS code [8]. Measuring the mode size from Fig. 3, we obtain

$$\sigma_r \cong 10 \mu\text{m}. \quad (18)$$

Note that this number agrees well with the above  $(\sigma_r)_{PAL}$  obtained by the formula of Eq. (14).

## CONCLUSION

Comparison with the simulation shows that the derived formula Eq. (14) is fairly accurate.

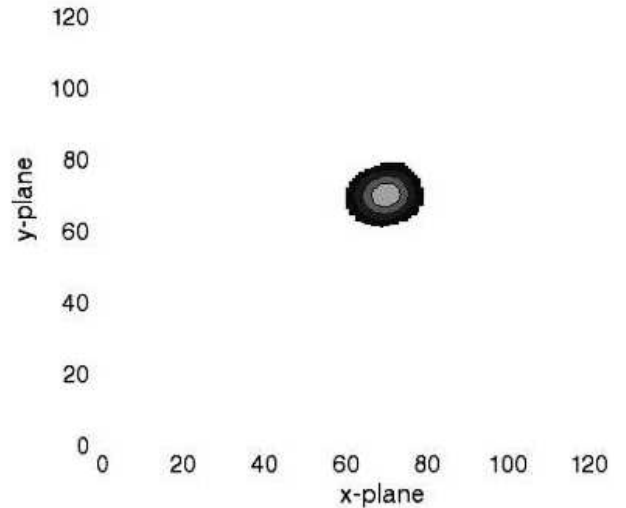


Figure 3: Simulated transverse mode size of the PAL-XFEL.

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