

# COLLIMATION OPTIMIZATIONS, CAPTURE EFFICIENCY, AND PRIMARY-BEAM POWER LOSS IN THE ILC POSITRON SOURCE\*

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## Abstract

The ILC positron beam generated from a thin Ti target has a wide energy spread and large transverse divergence. With the collection optics immediately downstream of the target and pre-acceleration to 125 MeV, the collected positron beam still has a long tail of positrons with low energies and large transverse divergence, which will be lost in the rest of the ILC positron source beamline. A collimation system is proposed and optimized for the case of a shielded target with quarter-wave transformation collection optics so that the power loss in the magnets and RF structures is effectively controlled within the acceptable level and in the damping ring (DR) within 640 W, assuming  $3 \times 10^{10}$  of the captured positrons per bunch in the DR. In this case, the capture efficiency and DR injection efficiency are 13% and 99.8%, respectively. The lower capture efficiency is expected to result in higher injection efficiency and therefore, a lower power loss in the DR. The capture efficiency for the cases of a shielded target with flux concentrator and 5-T immersed target with flux concentrator is 20% and 30%, respectively, with the collimation system.

## OVERVIEW OF THE ILC POSITRON SOURCE OPTICS

The 150 GeV International Linear Collider (ILC) electron beam passing through an undulator generates circularly polarized photons, which impinge on a thin Ti target and produce longitudinally polarized positrons. The generated positron beam is first collected and accelerated to 125 MeV through a beamline TAP. Then a dogleg, PCAP, is used to separate positrons from electrons and photons. Positrons are accelerated to 400 MeV in the normal conducting (N.C.) pre-acceleration linac, PPA. The 400 MeV beam passes through a beamline, PPATEL, to the electron main linac tunnel. Then the positron beam transports through a beamline, PTRAN, from the electron main linac tunnel to the positron superconducting (S.C.) booster linac, PBSTR. After acceleration to 5 GeV, it is transported from the PBSTR Linac-to-Ring (LTR) performing spin rotations and energy compression, and finally enters the damping ring (DR) injection line. The optics and geometry are described in detail in Refs. [1-2].

## OPTIMIZED COLLIMATION SYSTEM

A collimation system is designed to reduce the beam power loss in the beamline elements (e.g. magnets, RF

structures, and drift beam tube) to an acceptable level, typically smaller than 100 W/m and 1 W/m for N.C. and S.C. elements, respectively, without a severe loss of capture efficiency of the positron source transport. Here, we define the capture efficiency as the number of positrons captured inside the DR acceptance divided by the initial number of positrons at the target. The positron beam entering the transport line has a long tail characterized by low energy and large divergence. Without the beam collimation, the beam loss in the beamline elements would severely exceed the acceptable level, particularly in the first three sections: the PCAP, PPA, and PPATEL system. Following the preliminary design of collimation system in [3], total of eleven collimators are proposed in the PCAP, PPA, and PPATEL in this design. The first four collimators in the beginning of PCAP are used for betatron amplitude collimation of the incoming positrons with large transverse amplitudes and angles, while the next four collimators positioned in the dispersive areas are mainly used to clip off the low energy tail. The combination of the ninth and tenth collimators is to effectively reduce the beam loss in the N.C. RF structures in the PPA immediately downstream of the PCAP. The eleventh collimator placed at the dispersive location in the PPATEL is to decrease the beam power loss in the magnets below 100 W/m. In principle, the final collimation of the positrons which are beyond the DR transverse acceptance can be done in the 5-GeV LTR section immediately upstream of the DR, but due to the high beam energy the power load on the LTR collimators would be relatively high. Instead, the collimation of the positrons beyond the DR normalized transverse acceptance,  $\gamma(A_x + A_y) \leq 0.09$  m, is performed

in the lower energy region,  $\leq 400$  MeV, by extensively optimizing the apertures of the first eleven collimators. With the current optics configuration, the collimation of the positrons beyond the DR longitudinal acceptance,  $\Delta E \times \Delta z \leq (\pm 25 \text{ MeV}) \times (\pm 3.46 \text{ cm})$ , has to be done at the LTR. Total of five energy collimators indexed from 12<sup>th</sup> to 16<sup>th</sup> are used for this purpose. Parameters of the complete collimation system along the positron transport are shown in the Table I. All collimators are assumed rectangular. The details of the collimation system are described in [4].

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Table I. Parameters of the rectangular collimators.

Collimator index	Half aperture in x/y (mm)	Length (cm)	Entrance location s (m)
In PCAP beamline			
C1	15/15	10	0.2
C2	28/25	10	1.0
C3	22/22.5	10	2.6
C4	38/25	10	4.27
C5	23/75	10	10.321
C6	25/75	10	11.570
C7	30/60	10	12.320
C8	42/18	10	15.319
C9	23/30	10	61.672
C10	16/16	10	64.792
In PPATEL: C11	15/7.5	8	109.292
In LTR beamline			
C12	23/35	20	5501.887
C13	30/35	30	5579.153
C14	20/35	60	5579.756
C15	7.4/35	50	5583.350
C16	10.5/35	60	5587.643

## PRIMARY-BEAM POWER LOSS AND CAPTURE EFFICIENCY

Primary-beam tracking from a thin Ti target to the entrance of the DR injection line has been performed. The tracking from the target to the capture section (125 MeV) is described in Ref. [5]. The Elegant code [6] is used to track the positron beam through the rest of the beamline including the PCAP, PPA, PPATEL, PTRAN, PBSTR, and finally the LTR system. Positron 6-D coordinates at the exit of the capture section are used as the input data for the Elegant code tracking. Note that only the positrons in the main RF bucket are selected for the tracking in this study. Due to the extremely large energy spread in the beginning of the beamline, the tracking was set up to calculate energy dependence to all orders in the magnets from PCAP to PTRAN sections, and then to the 2<sup>nd</sup> order for the rest of the transport, where the energy spread is reduced. To maximize the number of positrons within the DR acceptance, the energy compression in the LTR is fully optimized before the positrons reach the DR injection line. For that purpose, the booster linac PBSTR upstream of the LTR runs its RF phase off-crest to create a suitable correlated energy spread. The collimation system described in the previous section is implemented and the physical apertures of the beamline listed in Table II are used in the tracking.

The full beam power is 320 kW at the 5-GeV DR based on the ILC beam parameters –  $3 \times 10^{10}$  of captured positrons in the DR per bunch (50% more than the design value at the IP), 2670 bunches per pulse, and 5 Hz pulse repetition. The primary-beam power loss along the beamline is shown in Fig. 1 for the case of a shielded target with quarter-wave transformation optics. It is shown that the significant power loss is dissipated in the

sixteen collimators, and the power loss in N.C. components and S.C. booster linac is within 100 W/m and 1 W/m, respectively. The number of positrons per bunch along the transport is calculated, as shown in Fig. 2, which is used for the power loss calculation. The power loss in the five LTR collimators is about 5.5 kW, 10.7 kW, 13 kW, 15.6 kW, and 5 kW, respectively, and in the DR it is 0.64 kW. 13% of the positrons from the target survive the transport through the complete beamline, and 99.8% of the injected positrons are captured within the DR 6-D acceptance, corresponding to 12.97% of positrons from the target captured in the DR. Without the collimation system, the capture efficiency is 16.8%. Thus, 3.8% of capture efficiency is lost when the collimation system is used. The higher injection efficiency (i.e., the number of positrons captured in the DR divided by the injected positrons in the DR) corresponding to a lower power loss in the DR is expected to result in lower capture efficiency, as shown in Fig. 3. Tracking for various kinds of scenarios, such as immersed vs shielded target, flux concentrator vs quarter-wave transformation, is extensively conducted. The capture efficiency for the different field on the target and collection optics with and without the collimation system is shown in Fig. 4. It shows that the capture efficiency for the cases of a shielded target with flux concentrator and 5-T immersed target with flux concentrator is 20% and 30%, respectively.

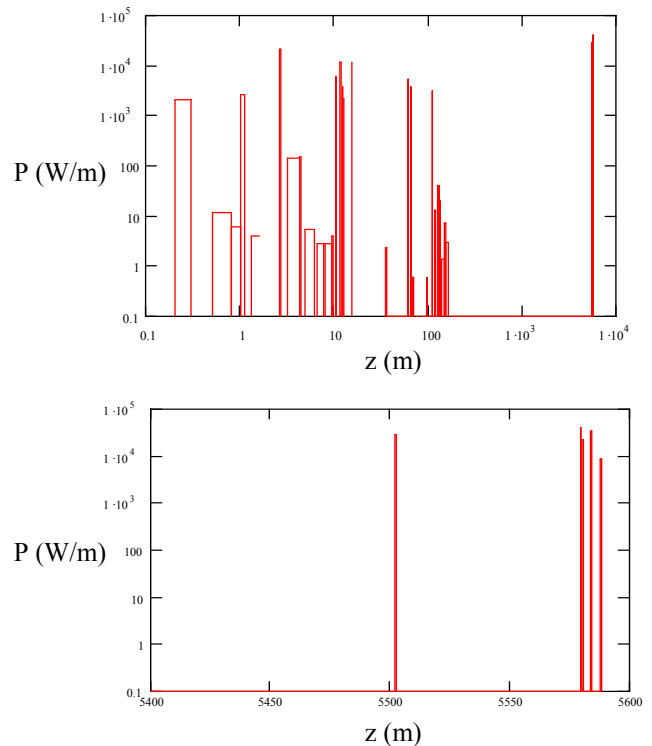


Figure 1: Primary-beam power loss along the complete positron source transport (top) and the LTR (bottom) given  $3 \times 10^{10}$  of the captured positrons in the DR; total of 16 collimators are implemented which absorb most of the power loss.

Table II. Physical apertures of the beamline.

Components	Half aperture in x/y (mm)
Capture section	23/23
PCAP	75/75
PPA	23/23
PPATEL	75/75
PTRAN	75/75
PBSTR	37/37
LTR	
RF section	37/37
solenoid	20/20
others	75/35

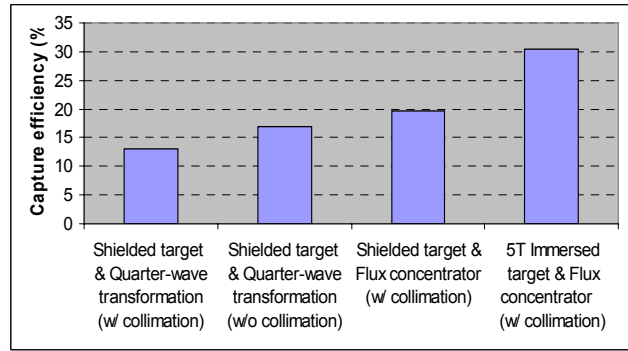


Figure 4: Capture efficiency for different field on the target and collection optics with and without the collimation system optimized for the case of a shielded target with quarter-wave transformation optics.

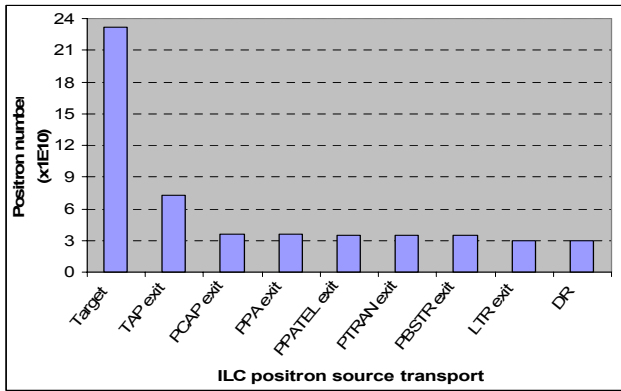


Figure 2: Number of positrons per bunch from the target to the DR.

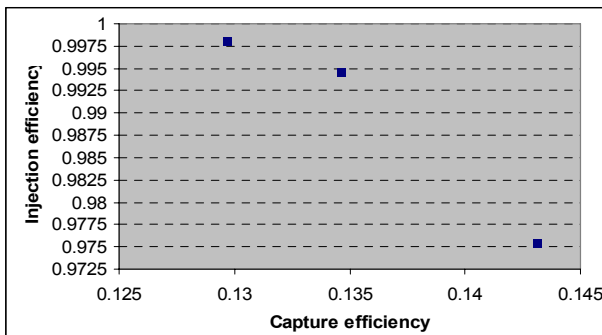


Figure 3: Injection efficiency vs capture efficiency; the lower capture efficiency is expected to have higher injection efficiency, resulting in a lower power loss in the DR.

### SUMMARY

A collimation system for the positron source transport is optimized for the case of a shielded target with quarter wave transformation collection optics. The primary-beam tracking shows that with the collimation system the beam power loss in the transport line can be controlled within the acceptable level and the power loss in the DR is 640 W for  $3 \times 10^{10}$  of captured positrons per bunch in the DR. The corresponding injection efficiency and the capture efficiency is 99.8% and 13%, respectively. The lower capture efficiency is expected to result in higher injection efficiency and therefore, a lower power loss in the DR. The capture efficiency for the cases of a shielded target with flux concentrator optics and a 5-T immersed target with flux concentrator is 20% and 30%, respectively, with the collimation system. Detail studies including the secondary particles were presented in Ref. [7]. We would like to thank Drs. V. Bharadwaj and W. Gai for helpful discussion.

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