DESIGN OF THE MULTI-ION INJECTOR LINAC FOR THE JLAB EIC (JLEIC)*

B. Mustapha, Z. Conway and M. Kelly, Argonne National Laboratory, Lemont, IL, USA A. Plastun and P. Ostroumov, Facility for Rare Isotope Beams, East Lansing, MI, USA

Abstract

An Electron Ion Collider (EIC) is the highest priority for future U.S. accelerator-based nuclear physics facility following the completion of the Facility for Rare Isotope Beams (FRIB). Two laboratories are competing to host the future EIC: Brookhaven National Lab. (BNL) and Jefferson Lab. (JLab). The baseline design of JLab's Electron Ion Collider (JLEIC) ion complex comprises a pulsed superconducting (SC) linac injector capable of accelerating all ions from protons to lead, where proton and light ion beams can be polarized. After reviewing the design requirements for the injector linac, important design choices such as the room-temperature (RT) section design, the transition energy between the RT and SC sections and the stripping energy for heavy ions will be discussed. The design of the different linac sections will be presented as well as the results of end-to-end beam dynamics simulations for polarized deuterons and un-polarized lead ions.

INTRODUCTION

The electron-ion collider concept proposed by JLAB (JLEIC) requires a new ion accelerator complex which includes a multi-ion linac capable of delivering any ion beam from hydrogen to lead to the Booster. We have developed a design for a pulsed linac which consists of different ion sources, a room-temperature (RT) front-end, up to 5 MeV/u followed by a superconducting (SRF) section to the full linac energy. This work includes the beam dynamics and electrodynamics studies performed to design efficient and cost-effective accelerating structures for both the RT and SRF sections of the linac. The current design includes two separate RFOs one for heavy ions and one for polarized light-ion beams, and a common RT section with a special IH DTL design downstream of the RFQs. Quarterwave and half-wave resonators are effectively used in the SRF section of the linac.

DESIGN REQUIREMENTS AND CHOICES FOR JLEIC INJECTOR LINAC

Design Requirements for JLEIC Injector Linac

The baseline design of the JLEIC ion complex [1] calls for the following requirements from the injector linac:

- Capable of accelerating all beams from protons to lead ions, including polarized light ion beams
- Deliver 280 MeV protons and 100 MeV/u lead ions Pb⁶⁷⁺ for injection to the Booster, and equivalent energies for other ion beams

- Pulsed beam structure with 5-10 Hz repetition rate and 0.2 0.5 ms beam pulse length
- Pulsed beam current of ~ 2 mA for light ions and ~ 0.
 5 mA for heavy ions
- Compact and cost efficient

Important Design Choices for the Linac

In order to satisfy the design requirements listed above, the following design choices were made for the JLEIC injector linac:

- To accommodate the significantly different beam parameters from polarized light-ion and heavy-ion sources, the linac includes two separate RFQs, one for mass-to-charge ratio $A/q \le 2$ and one for heavy ions with A/q > 2.
- As a consequence, two separate low-energy beam transport (LEBT) lines are required. However, this separate front-end choice allows a special LEBT design for polarized light ions to preserve polarization.
- Based on similar pulsed ion linacs [2, 3], a roomtemperature (RT) section up to an energy of ~ 5 MeV/u is the most efficient and cost-effective option for the JLEIC linac, followed by a SRF section up to the full linac energy.
- A pulsed SRF linac can be more compact and costeffective than the full RT option [4, 5]. It also offers wider acceptance and more tuning flexibility for light and heavy ion beams. In addition, taking advantage of state-of-the-art performance of quarter-wave (QWR) and half-wave (HWR) resonators [6, 7], which can deliver higher voltages in pulsed mode, the linac can be even more compact.
- In order to deliver Pb⁶⁷⁺ at 100 MeV/u, the optimum stripping energy was found to be ~ 13 MeV/u, which is the energy following two QWRs modules made of 7 cavities each.

DESIGN OF THE DIFFERENT SECTIONS OF THE LINAC

Figure 1 shows the layout of the designed JLEIC injector linac with separate front-ends for light ion and heavy-ion beams, a DTL section made of three IH tanks followed by an SRF section made of three QWR cryomodules operating at 100 MHz and nine HWR cryomodules operating at 200 MHz. A stripper section for the heaviest ions is located between the second and third QWR modules.

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Figure 1: Layout of the JLEIC multi-ion injector linac.

We should mention that the current design is an extension of the short linac design [8] to the full energy required for injection to the booster ring in the JLEIC baseline design [1]. The short linac was proposed as part of an alternative design approach for the JLEIC Ion complex [9].

Existing Ion Sources and Expected Performance

Table 1 shows the desired polarized proton and deuteron ion beam parameters in comparison to the beams available from existing polarized sources, namely the atomic beam (ABPIS) and the optically pumped (OPPIS) sources [10].

Table 1: Polarized Light Ion Beams: Desired vs. Available H-/D- Beam Parameters

| | Units | Desired value | ABPIS value | OPPIS value |
|---------------|-------|------------------|----------------|----------------|
| Pulse Current | mA | 2 | 3.8 | 4 |
| Pulse Length | ms | 0.5 | 0.17 | 0.3 |
| Polarization | % | 100 | 90 | 85 |

We notice that while the desired beam parameters are not exactly matched, they are within reach with some R&D effort. For heavy-ion beams, both the electron cyclotron resonance (ECR) and the electron beam (EBIS) ion sources are capable of delivering the desired beam current of 0.5 mA for Pb³⁰⁺ and other ions with equivalent mass-to-charge ratios.

One notable difference between the polarized light ions and heavy ions is the beam emittance at the source, while the 90% emittance is typically 0.5π mm.mrad for heavy ions, it's ~ 2π mm.mrad for polarized H-/D- beams [11]. This significant difference in beam emittance is the main reason to have two separate front-ends for polarized lightion and heavy-ion beams.

LEBTs Design

The design for the polarized light ion LEBT shown in Figure 2 is very similar to the BNL LEBT for polarized H-[12]. It includes two opposite bends to compensate beam polarization. The LEBT is designed for 20 keV/u beams which is the RFQ injection energy.





The heavy-ion LEBT is shown in Figure 3, it is based on the CERN Lead linac-3 injector design [13]. The double-bend is for the selection of the desired charge state for injection into the RFQ. The extraction beam energy from the source is 25 keV/u which is also the RFQ injection energy.



Figure 3: LEBT design for heavy ion beams.

RFQs Design

The design parameters for the light ion and the heavy ion RFQs are listed side by side in table 2. Note the ~ 100 % transmission was achieved for a deuteron beam in order to avoid radio-activation of the RFQ structure by neutrons [14].

Table 2: Design Parameters for the Light-ion and Heavyion RFQs

| Parameter | Units | Light Ion | Heavy Ion |
|--------------------------------|-----------|--------------|--------------|
| Voltage | kV | 103 | 70 |
| Average radius | mm | 7.3 | 3.4 |
| Length | m | 3 | 5.6 |
| Quality factor | | 7200 | 6600 |
| RF power | kW | 150 | 250 |
| Beam transmission | % | ~ 100 | 99 |
| Design Transverse emittance | π.mm.mrad | 2.0 | 0.5 |
| Output Longitudinal emittance | π.keV.ns | 5.0 | 4.5 |

For pulsed RFQs, there are different options for the structure design, namely the 4-rod, the 4-vane and the window-coupled design as shown in Figure 4. The

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structure can be either brazed or bolted. For the JLEIC RFQs, we propose the brazed 4-vane window-coupled structure for both its mechanical and field stabilities [15].



Figure 4: RFQ Design options.

DTL Section Design

During the design of the RT section [16] of the JLEIC linac, we investigated different DTL design options, see figure 5. The first was an IH structure with triplet focusing similar to the BNL EBIS injector [2]. The second option was a special DTL design with RF quadrupole focusing [17]. The third option, which was selected for the JLEIC linac, uses a FODO focusing lattice which offers a large acceptance with minimal emittance growth while preserving a good power efficiency.





The design parameters for the three DTL tanks required to reach 5 MeV/u energy for all ion beams are listed in Table 3. The DTL section delivers \sim 30 MV over \sim 11 m and requiring ~ 1 MW total power. A schematic of the DTL section showing the three tanks made of 19 accelerating gaps and 20 quadrupoles in a FODO arrangement is shown in Fig. 6.

Table 3: Design Parameters for the Three IH-DTL Tanks with FODO Focusing

| Parameter | Units | DTL-1 | DTL-2 | DTL-3 |
|-----------------|-------|-------|-------|-------|
| Input Energy | MeV/u | 0.5 | 2.0 | 3.6 |
| Output En. | MeV/u | 2.0 | 3.6 | 5.0 |
| Accel. sections | No. | 10 | 5 | 4 |
| Length | m | 4.3 | 3.5 | 3.4 |
| RF power | kW | 280 | 400 | 620 |
| | | | | |
| IH-1 | | IH-2 | I | H-3 |
| | | • • • | | |

Figure 6: Layout of the DTL section made of 3 tanks including 19 accelerating gaps and 20 quadrupoles in a FODO lattice arrangement.

shown in Fig. 7. It is made of three QWR and nine HWR cryomodules. Each cryomodule is made of seven cavities and four solenoids in the arrangement shown in Figure 8 for both the QWR and HWR modules. A stripping section for the heaviest ions is located after the second OWR

SRF Section Design



module at an energy of ~ 13 MeV/u for lead ions.

Figure 7: Schematic layout of the SRF section of the JLEIC including a stripper section after the second QWR module.



Figure 8: Lavouts of the OWR and HWR cryomodules each with seven cavities and four solenoids.

The design of both the QWR and HWR cavities are shown in Fig. 9 along with their electromagnetic field distributions. The corresponding RF design parameters are summarized in Table 4.

QWR E&M Fields

HWR E&M Fields



Figure 9: Geometry and electromagnetic field distributions for the QWR and HWR cavities (E-Field on the left and B-Field on the right for each cavity).

Table 4: RF Design Parameters of the QWR and HWR Cavities

| Parameter | Units | QWR | HWR |
|----------------------------------|-----------|-------|-------|
| Design β_{opt} | - | 0.15 | 0.3 |
| Frequency | MHz | 100 | 200 |
| Length ($\beta\lambda$) | cm | 45 | 45 |
| $E_{\text{peak}}/E_{\text{acc}}$ | - | 5.5 | 4.9 |
| $B_{\text{peak}}/E_{\text{acc}}$ | mT/(MV/m) | 8.2 | 6.9 |
| R/Q | Ω | 475 | 256 |
| G-factor | Ω | 42 | 84 |
| E _{peak} in operation | MV/m | 58 | 52 |
| B _{peak} in operation | mT | 86 | 73 |
| Eacc | MV/m | 10.5 | 10.5 |
| Voltage per cavity | MV | 4.7 | 4.7 |
| Cavity phases | Deg | 15-30 | 15-30 |
| No. of Cavities | - | 21 | 63 |

Stripper Section Design

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The Pb beam is used as a reference for the design of the stripper section. The optimum stripping energy on a carbon foil to produce lead ions Pb⁶⁷⁺ for injection to the JLEIC booster is about 13 MeV/u [18]. This energy maximizes the beam fraction in the desired Pb⁶⁷⁺ charge state which is about 20% and minimizes the total voltage requirements for the linac up to the full-energy of 100 MeV/u for lead ions as shown in Fig. 10.



must maintain attribution to the author(s), title Figure 10: Total linac voltage as function of stripping energy for lead ions.

work In order to separate the desired charge state from unthis wanted charge states and other reaction products, a chicane of can be used for a straight linac option or a 180-deg bend distribution for a folded option. Figure 11 shows a preliminary concept of a stripping chicane, where the beam is focused onto the stripper foil using a triplet and the desired charge state is separated and selected using the slits in the middle plane. Anv A rebuncher and another triplet are used at the end of the chicane to longitudinally and transversely match the beam BY 3.0 licence (© 2018). to the following section.



Figure 11: A possible layout for a stripping chicane to be under the terms of the CC installed between the second and third QWR cryomodules.

BEAM DYNAMICS IN THE LINAC

Beam Dynamics Design

The beam dynamics design is straightforward for the SRF section of the linac, a focusing period made of two accelerating cavities and one solenoid is the main building block. Exception is made at the end of every cryomodule where a missing cavity accounts for the extra drift space between modules where beam diagnostic devices can be installed. The general design rule is to start with a phase work may advance below 90-deg for the zero current beam and maintain periodic focusing by smoothly decreasing the phase advance along the linac. The accelerating voltage profile in the cavities is shown in Fig. 12. It clearly shows that the HWR covers very well the velocity range from 0.2 to 0.35 and no need for a different cavity type at the higher energies. This simplifies the overall design and fabrication of

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Figure 12: Effective cavity voltage as function of β along the linac showing that the HWR is fairly efficient up to the full energy of the linac.

Beam Dynamics Simulations

The results of end-to-end beam dynamics simulation in TRACK for a 2-mA polarized deuteron beam are shown in Fig. 13. The simulation starts from the ion source through the RT section, and ends with the SRF section at the full linac energy required for injection to the booster. In this case, we note $\sim 100\%$ transmission of the deuteron beam with about 30% emittance growth. Similar results were obtained for a 2-mA polarized proton beam.



Figure 13: End-to-end beam dynamics simulation results for deuteron beam with $\sim 100\%$ transmission through the whole linac.

For lead ions, the beam dynamics simulation results are shown in Fig. 14.

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Figure 14: End-to-end beam dynamics simulation results for lead ions with beam stripping at ~ 13 MeV/u.

SUMMARY

We have designed a pulsed multi-ion injector linac that satisfies all the requirements for the baseline design of the Jefferson Lab EIC concept (JLEIC). The design is based on a 5 MeV/u room-temperature section followed by an SRF linac using QWRs and HWRs taking advantage of the most recent developments of similar structures at Argonne.

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REFERENCES

- Y. Zhang *et al*, "JLEIC Ultimate Luminosity with Strong Electron Cooling", in *Proc. of IPAC17*, Copenhagen, Danmark, May 14-19, 2017
- [2] A. Schempp *et al*, "RFQ and IH accelerators for the new EBIS injector at BNL", in *Proc. of PAC2007*, Albuquerque, New Mexico, USA
- [3] N. Angert *et al*, "The IH Linac of the CERN Lead Injector", in Proc. LINAC-1994, Tsukuba, Japan
- [4] P. Ostroumov, "Warm and Cold Ion Linac: Comparison and Optimization", Presentation at the MEIC Collaboration Meeting, JLab, Spring 2015
- [5] R. York, "Status Report on Technology Evaluation for JLEIC Ion Linac", Presentation at the JLEIC Collaboration Meeting, JLab, Spring 2016
- [6] M. Kelly *et al*, "Low Beta Cavity Development for an AT-LAS Intensity Upgrade", in *Proc. SRF2013*, Paris, France
- [7] Z. Conway *et al*, "Achieving High Peak Fields and Low Residual Resistance In Half-Wave Cavities", in *Proc. SRF2015*, Whistler, BC, Canada

Table 5 summarizes the beam parameters from the linac for protons, deuterons and lead ions.

Table 5: Important Beam Parameters in the Linac for Pro tons, Deuterons and Lead Ions

| Parameter | H- | D- | Pb ³⁰⁺ → ⁶⁷⁺ |
|---|------|------|------------------------------------|
| Input Transverse emit- tance (π .mm.mrad) | 2.0 | 2.0 | 0.5 |
| Output energy (MeV/u) | 299 | 172 | 110 |
| Transmission (%) | 99.7 | 99.9 | 97.9 |
| Output Transverse emitance (π.mm.mrad) | 3.0 | 3.0 | 1.0 |
| Output Longitudinal emittance (π.keV/u.ns) | 7.5 | 7.5 | 7.0 |
| Beam energy spread (%) | 0.1 | 0.1 | 0.1 |

- [8] P. Ostroumov *et al.*, "Design and Beam Dynamics Studies of a Multi-ion Linac Injector for the JLEIC Ion Complex", in *Proc. of HB2016*, Malmo, Sweden, July 2016, paper THPM5Y01, p. 559.
- [9] B. Mustapha et al., "An Alternative Approach for the JLEIC Ion Accelerator Complex", in Proc. NAPAC2016, October 9-14, Chicago, IL
- [10] A. Sy and V. Dudnikov, "JLEIC Ion Source: Specifications, Design and R&D Prospects", Presentation at the JLEIC Collaboration Meeting, JLab, Fall 2016
- [11] A. S. Belov, "Polarized ion source with nearly resonant charge-exchange plasma ionizer: parameters and possibility for improvements", Proceedings of the 15th International Workshop on Polarized Sources, Targets and Polarimetry, September 9-13, 2013, Charlottesville, VA, USA
- [12] J. Alessi *et al*, "Design of a 35 keV LEBT for the new High Intensity OPPIS at BNL", in *Proc. PAC1999*, New York, NY, USA
- [13] N. Angert *et al*, "CERN heavy-ion facility design report", CERN Report 93-01, 28 April 1993, Geneva, Switzerland
- [14] R. Ferdinand *et al*, "Commissioning of SPIRAL2 CW RFQ and Linac", in *Proc. IPAC-2017*, Copenhagen, Denmark
- [15] P. Ostroumov *et al*, "Development and beam test of a continuous wave radio frequency quadrupole accelerator", *Phys. Rev. ST Accel. Beams*, 15 (2012) 110101
- [16] A. Plastun *et al*, "Design of the room-temperature front-end for a multi-ion linac injector", in *Proc. NAPAC2016*, Chicago, IL, USA
- [17] A. Kolomiets and A. Plastun, "Spatially periodic radio-frequency quadrupole focusing linac", *Phys. Rev. ST Accel. Beams*, 18 (2015) 120101
- [18] A. Leon et al, "Charge state distributions of swift heavy ions behind various solid targets", Atomic Data and Nuclear Data Tables 69 (1998) 217-238