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Charge Measurements in SwissFEL and Results of an Absolute Charge Measurement Method

10th International Beam Instrumentation Conference (IBIC 2021) Remote 13-17 September 2021

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➢ Overview of the SwissFEL charge monitors

- Faraday Cup
- Integrating Current Transformer
- Wall Current Monitor
- Cavity Beam Position Monitor

Wall Current Monitor: experimental characterization and numerical modelling results

Cavity Beam Position Monitor for absolute charge measurements

Charge measurement campaign at SwissFEL

Procedure for absolute calibration of the SwissFEL charge monitors

➤Conclusions



➢ Beam energy: 6.2 GeV and 3.3 GeV

Beam charge: 10-200 pC @100Hz, 28 ns 2-bunch time structure

Emittance: 0.4/0.2 mm mrad

➢Bunch length: from a 3 ps up to a few fs

Photon wavelength: 0.1-0.7 nm and 0.7-7.0 nm



- Turbo ICT (Integrating-Current-Transformer):
 - 4 in Aramis and 2 in Athos
 - 28 ns 2-bunch resolution
- Standard ICT (1 in the gun section):
 - full beam current integral (no 2-bunch discrimination)
- Cavity-BPM (Beam-Position-Monitor):
 - ~200 units distributed all along Aramis and Athos
- > Wall-Current-Monitor (WCM, 1 in the gun section):
 - coarse synchronization photocathode laser timing and RF gun phase
 - 2-bunch charge monitoring @100 Hz
- Faraday- Cup (one unit in the Gun section):
 - Dark current



Standard ICT (BCM-IHR readout electronics)

- time integration (5µs) of the beam induced current in the transformer
- Broad-band frequency response (kHz→10 MHz)
- no 2-bunch time structure resolution
- Sensitive to dark current of the gun
- Charge resolution ~ pC



Courtesy Bergoz



SwissFEL ICT: dual-bunch resolving

> Turbo-ICT2 (BCM-RF readout electronics)

- High frequency transformer (bandwidth up to several hundred MHz)
- Narrow band-pass filter centered at around 180 MHz (resonance quality factor reduced to discriminate the 2 bunch)
- Output signal is a resonance with amplitude proportional to the bunch charge
- Beam charge determined by measuring the apex of the resonance (sample-and-hold electronics)
- 28 ns 2-bunch time structure discrimination and immunity to dark current
- rms resolution 0.1 pC (1%) in the 10-200 pC SwissFEL charge range





Faraday Cup (FC)

➤The Faraday cup destructively intercepts the beam

Low current and dark current measurements

Systematic error in charge measurements:

- "Containment": mismatch between e.m. shower and absorber dimensions
- "Charge trapping": absence of voltage o magnetic "cage" to bring back secondary and scattered electron
- resistor signal coupled with DC component of input current





M. Dach et al. (SLS Linac) ; BIW2000





Wall Current Monitor in SwissFEL



Prototype Wall Current Monitor In WLHA $R_{gap} = 3.0 + /- 0.05 \Omega (2.83\Omega with 50\Omega)$ (12 x 36 Ohm gap resistors) NiZn ferrite ring



$$Q = \frac{1}{R_{gap}} \int V_{Rgap}(t) dt$$

SwissFEL Wall Current Monitor

After RF gun (z=0.58 m) **R_{gap} = ?** NiZn ferrite ring



Time integration window ~ 10 ns

Lower cutoff frequency $f_{low}=R_{gap}/2\pi L \sim 200$ kHz (no DC component of input current measurable) Upper cutoff frequency $f_{high}=1/(2\pi R_{gap} C) \sim G$ Hz good WCM \rightarrow flat transfer impedance Z(ω) $\sim Rgap$ in the frequency range of interest (up to hundreds of MHz) PAUL SCHERRER INSTITUT

CST Simulation WCM output (25 ps , 200 pC)





time response

- The Gaussian beam signal is distorted due to the upper cutoff frequency
- Voltage oscillations in the resistor (long time regime)

frequency response

- Approx. constant for f << 1 GHz
- Broad peak at 1.5 GHz
- Sharp peak at 4.7 GHz (voltage oscillations)

1.5 Gz peak → mismatch of the gap resistance and transmission line capacitive gap

4.7 GHz sharp peak \rightarrow step transition (radius) from the beam pipe to the gap

Charge Measurement in the simulation



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- The measured charge decreases.
 - No DC measurable. For infinite integration time, Q=0.
- The slope is determined by the lower cutoff frequency $f_{low} = \frac{R_{gap}}{2\pi L}$
- First few data points after bunch incidence include current components outside the WCM bandwidth (non measurable)

Idea: Make linear fit of slope and intersect fit with the step to reconstruct the charge at the time when the bunch arrives.

- Holdoff (2.5 ns) to avoid first few data points
- Fitted interval = 10 ns



Charge Measurement: the Ferrite Ring



- Dispersive ferrite yields charge losses, because we only measure the voltage over the gap
- > Which value of the gap resistance to be used for WCM charge measurement ?



SwissFEL WCM: gap resistance value?

Spare WCM prototype \rightarrow Multimeter inspection and VNA measurements of the reflection coefficient S11 at the output port confirmed a gap resistance Rgap=3.0 Ω .

SwissFEL WCM \rightarrow No inspection possible, VNA measurements of S11 \rightarrow Rgap=3.55 Ω .

- > Transfer impedance measurements of spare WCM in a test-bench (coaxial wire method):
 - Transfer impedance $Zt(\omega)$ constant up to several MHz \rightarrow $Z(\omega)$ ~Rgap
 - Confirmed reduction of 5-6% of the transfer impedance because of ohmic losses in the ferrite Zt(ω) = 2.68 Ω instead of 2.83 Ω.
 - estimate of the ohmic resistance of the ferrite ~50.50 Ω





$$Z(\omega) = \frac{V_{out}(\omega)}{I_{beam}(\omega)} [\Omega]$$





SwissFEL WCM transfer impedance $Zt(\omega)$?

Direct measurement of transfer function of the SwissFEL WCM not possible. Possible Solution:

- Replace the WCM currently installed at SwissFEL with the prototype WCM which is now fully characterized.
- Get Z_t from a CST simulation under hypothesis of two resistors disconnected and dispersive ferrite (Simulation parameters: Rgap = 3.31 Ω, 25 ps bunch length, Q = 200 pC)
- > to compensate ohmic losses in ferrite instead of Rgap = $1/(1/3.55+1/50) = 3.31 \Omega$, the best up to date estimate of the transfer impedance are:
- Rt \approx 3.26 Ω CST prediction (near-DC value of the curve)
- Rt \approx 3.11 Ω (spare WCM VNA measurements)
- CST prediction sensitive to the permeability model of the ferrite
- To date, optimistic estimate of the incertitude on the WCM calibration parameter (Rt) about +/-5%



SwissFEL WCM: waveform signal integration and charge results

200 pC (charge readout of first turbo ICT)

 $R_t = 3.26 \ \Omega$ (derived from SwissFEL WCM Zt simulation, considering dispersive ferrite)



We integrate a background signal. Slopes are approximately linear.

Still possible to measure the charge with our method using a linear fit! Adapted idea: We extrapolate the charge and background signal contribution back to the time where the bunch was incident.



Cavity Beam position Monitor (CBPM)



- Two cavity device
- Reference cavity is designed to get an output signal proportional to the beam charge (monopole mode).
- Dependence on beam position is negligible (TM010 mode).
- > We used:
- CBPM16 type (low-Q).
- CBPM8 type (high-Q)







Cavity BPM and charge measurement

BPM output signal measured with a 16GHz, 40Gs/s oscilloscope.

Bandpass filter at the BPM output to isolate the cavity fundamental mode (TM010) signal.

BPM type	TM010 frequency	Low-pass filter cutoff frequency
CBPM16	3.284 GHz	4.8 GHz
CBPM8	4.926 GHz	6.0 GHz

Voltage induced in the TM010 cavity mode and available at the cavity output port is:

$$V_{out}(t) = q\omega \sqrt{\frac{Z}{Q_e}R/Q}e^{-\frac{\omega t}{2Q_L}}\cos(\omega t)$$

(e.g. see : Cavity Beam Position Monitors, R. Lorenz)

q(charge); Z(50 Ω , impedance cavity output line); ω (frequency of cavity mode TM010); R/Q (parameter depending on cavity geometry); Qe (external quality factor); QL (loaded quality factor)

For every single BPM ω , Qe and QL have been measured with VNA. R/Q is estimated with reliable numerical codes (HFSS and CST, same result).

To better reproduce the measured signal, the expression above is also convolved with a low-pass filter function, same cutoff frequency as the filter used for measurements.



high Q (left) and low Q (right) BPM: waveforms measured with the oscilloscope and fit results



Beam charge measurements at SwissFEL





With respect to first ICT (~200 pC) about 20pC smaller mean charge readout from cavity BPMs:

Left (B1) → (10+/-3)%

Right (B1)→ (10+/-2)%

Left: Charge measurements (B1, B2) of cavity BPMs and closest ICT in ARAMIS and ATHOS (measurements not simultaneous)

Right: charge measurements (B1) of cavity BPMs in high energy part of ARAMIS compared with average value of charge readout from the first Turbo-ICT-2 at the gun



Beam charge measurements at SwissFEL

Simultaneous charge readouts of first cavity BPM and at the gun and the WCM for variable charge and three different gun settings: (1) Only bunch1; measured bunch 1 in presence of bunch-2; measured bunch-2 in presence of bunch-1.

Nominal charge	200pC	100 pC	10 pC
WCM/BPM bunch-2 only	1.0092	0.9773	1.0361
WCM/BPM bunch-1	1.0011	0.9770	1.0204
WCM/BPM bunch-2	0.9953	0.9819	1.0341

WCM transfer impedance

estimate Rt \approx 3.26 Ω (CST prediction)



In the case bunch-1 at 200 pC, the relative percentage difference of the charge readout of the 1st cavity BPM w.r.t. the 1st ICT is about:

(8+/-1)%



Conclusions and Outlook

≻WCM:

- experimental characterization (VNA) and numerical modelling (CST, HFSS) but still incertitude on the calibration (+/-5%, very optimistic).
- Further characterization needed to solve the incertitude
- Main problem: no VNA transfer function measurements possible in the SwissFEL WCM

Cavity BPM:

- New method developed and implemented at SwissFEL for absolute charge measurements
- Robust and statistically consistent reliability for charge measurements at SwissFEL
- The aggregate results of the campaign of measurements with cavity BPM give us a calibration reference for alignment of all the charge monitors in SwissFEL
- Calibration procedure of charge monitors for bunch-1: apply a correction factor (9+/-2)% to the charge readout measured by the first Turbo-ICT-2 at the gun and align all the other charge monitors under a condition of full transmission along the entire machine