# LATTICIES FOR MILLI-eV NEUTRAL MOLECULES \*

H.Nishimura, G.Lambertson, J.G.Kalnins, and H. Gould, LBNL, Berkeley, CA94720, USA

#### Abstract

An electric dipole moment of neutral polar molecules interacts with non-uniform electric field; therefore it is possible to design electrodes that replace magnets for charged particles and a lattice that is a series of electrodes to circulate such molecules. We describe our recent result of designing a ring for  $CH_3F$  in comparison with our previous design for  $ND_3$ .

## STARK POTENTIAL

The interaction between the dipole moment of a neutral polar molecule and an external electric field is defined by the Stark potential W. As a dipole always lines up to the field direction, the potential W is a function of the field strength  $E = |\vec{E}|$  of the external field  $\vec{E}$ . The x force is given by:  $F_x = -\partial W/\partial x = -dW(E)/dE \cdot \partial E/\partial x$  and similarly for y and z. The molecule is called strong-seeking if dW(E)/dE > 0, or weak-seeking if dW(E)/dE < 0. ND<sub>3</sub> is an example of the weak-seeking case with the following potential:

$$W = \sqrt{C_1 + C_2 E^2} - C_3 - C_4 E^2 \tag{1}$$

where  $E = |\vec{E}|$ ,  $C_1 = 2.77 \cdot 10^{-49}$ ,  $C_2 = 6.35 \cdot 10^{-60}$ ,  $C_3 = 5.26 \cdot 10^{-25}$  and  $C_4 = 1.78 \cdot 10^{-38}$ .

CH<sub>3</sub>F is an example of the strong-seeking case with the following potential:

$$W = \frac{C_1 w_a^2 R}{1 + C_2 w_a} \tag{2}$$

where  $w_a = d_e E/R$ ,  $d_e = 1.86 \times 3.36 \cdot 10^{-30}$ ,  $R = 0.85 \times 1.99 \cdot 10^{-23}$ ,  $C_1 = -0.2085$  and  $C_2 = 0.2445$ .

#### **ELECTRODE**

There are 3 kinds of static electrodes shown in Fig.1. The type 1 and 2 electrodes are lenses to focus or defocus the beam in the transverse planes. For weak-seeking molecules, Type 1 is horizontally defocusing, 2 is focusing, and 3 deflects towards the right. For strong-seeking molecules, the forces are reversed. The fields from these electrodes are gradients of the following potential:

$$\Phi(x,y) = -E_0(y + A_2y + A_3(x^2y - \frac{1}{3}y^3) + A_5(x^4y - 2x^2y^3 + \frac{1}{5}y^5))$$
(3)

In Type 3,  $E_0A_2$  determines the bending radius and  $A_3$  becomes a knob to adjust the focusing forces. In arc sections we use cylindrical geometry for calculations and the potential differs slightly from Eq.3.[2]. There is also focusing in



the direction of the the fringe field at each longitudinal end of an electrode. It is approximated for small excursions, numerically integrated and replaced by a thin lens on or near the edge. It acts as a vertically focusing lens for weakseeking molecules and defocusing for strong-seeking ones. See [2] and [3] for details.

### **STORAGE RING**

### **Design** Principles

**Requirements** We require the lattice to provide the following.

(1)Dispersion-free long straight sections for injection, RF bunchers and beam experiments.

(2)Adjustable betatron tunes and beta functions.

(3)Tolerance for the effect of the gravity force.

**Solutions** Solutions we have adopted are:

(1)Achromatic arc: Adjusting the horizontal betatron phase advance in each arc to be a multiple of  $2\pi$ , the straight sections become dispersionless. The vertical phase advance is also tuned in the same manner to keep the vertical closed-orbit distortion(COD) due to the gravity in a reasonable range.

(2)A triplet of focusing/defocusing straight electrodes on each end of long straight sections to adjust betatron tunes and beta functions.

(3) Vertical orbit correction to compensate the effect of gravity if needed.

## Storage Ring for ND<sub>3</sub>

In case of a weak-seeking molecule, the bending electrode can be focusing in both horizontal and vertical planes. Therefore, it is possible to design a storage ring lattice with arc electrodes that bend the beam through a large angle [1]. The ring has the racetrack shape shown in Fig.2. Straight

<sup>\*</sup>Work on the synchrotron storage ring is supported by the Director, Office of Science, of the U.S. Department of Energy, and work on the linear decelerator is supported by the Director, Office of Science, Office of Basic Energy Sciences, of the U.S. Department of Energy; both under Contract No. DE-AC03-76SF00098.



Figure 2: Storage Ring Lattices

sections are 40 cm long. Acceptances are 35 mm  $\cdot$  mrad horizontal and 71 mm  $\cdot$  mrad vertical at its nominal velocity of 90 m/sec. Other ring parameters are listed in Table.1. Simple adjustments allow one to reduce the velocity to 60 m/sec [2].



Figure 3: Beta and Dispersion of the ND3 Ring

## Storage Ring for CH<sub>3</sub>F

For the strong seekers, a bending electrode can not focus in both planes simultaneously. Therefore, the arc must be an alternating sequence of focusing(BF) and defocusing(BD) bending electrodes; four cells with 90 ° phase advance make up a 45 ° arc sector with radius 60 cm. Parameters listed in Table.(2). As a result, the CH<sub>3</sub>F ring becomes much larger than the ND<sub>3</sub> ring as shown in Fig.2.



Figure 4: Beta and Dispersion of the CH3F Ring

The horizontal beta becomes large at Q2 which would cause strong amplitude-dependent tune shift. Therefore, the  $A_5$  component is introduced to Q2 to reduce the

Table 1: Ring Parameters							
Paramter	$ND_3$	CH <sub>3</sub> F					
Circumference (m)	3.357	9.850					
Circulation period (s)	0.0380	0.3121					
Velocity in free space (m/s)	90.0	30.0					
Symmetry of the ring	2	8					
Bending radius (m)	0.20	0.60					
Long Straight Section (m)	0.40	0.40					
Beta horizontal <sup>1</sup> : $\beta_x$ (m)	1.264	0.274					
Beta vertical <sup>1</sup> : $\beta_y$ (m)	0.513	0.596					
Horizontal dispersion <sup>1</sup> (m)	0.001	0.000					
Horizontal tune: $\nu_x$	5.250	13.368					
Vertical tune: $\nu_y$	5.200	10.398					
Chromaticity - H: $\zeta_x$	-0.0885	-73.8					
Chromaticity - V: $\zeta_y$	-0.0942	-38.1					
Momentum compaction: $\alpha$	-0.99	-0.899					
Dynamic aperture - H: $a_x$ (mm)	6.5	1.75					
Dynamic aperture - V: $a_y$ (mm)	6.0	3.50					
Acceptance - H: $\epsilon_x$ (mm - mr)	35	11					
Acceptance - V: $\epsilon_y$ (mm - mr)	71	21					

Table 2: Bending Electrode Parameters for  $CH_3F$  Ring, $E_0$  in MV/m, R and L in cm, T in degree.

	,		, U			
	Eo	$A_2$	$A_3$	R	T	L
BF	4.00	-10.55	-2296	60	5.625	7.85
BD	4.00	-10.55	2343	60	5.625	7.85

amplitude-dependent tune shift and increase the dynamic aperture.

The effect of gravity becomes an issue because of the lower velocity of the molecule. In case of the ND<sub>3</sub> ring, a small 0.17 mm displacement of the closed orbit did not require orbit correction. Note that in the CH<sub>3</sub>F ring, the sequence inside the triplets is changed to D-F-D as listed in Table 3. This reduces the effects of edge focusing but the vertical displacement still becomes as large as 2.6 mm and the dynamic aperture is completely killed. A vertical kick is needed, one choice is to shift Q2 downwards 0.24 mm to reduce the orbit distortion to 0.26 mm as shown in Fig.5. The dynamic aperture in Fig.6 is then unaffected by gravity.



Figure 5: Corrected Vertical COD of the CH3F Ring

A buncher is a pair of parallel plates with field pulsed in an offset triangular waveform that sweeps between zero



Figure 6: Dynamic Aperture of the CH3F Ring

Table 3: Parameters of Straight Electrodes for  $CH_3F$  Ring,  $E_0$  in MV/m and L in cm.

	Eo	$A_2$	$A_3$	$A_5$	L
Q1	3.34	0	2000	0	3.0
Q2	3.71	0	-2000	-1.283E+6	4.0
Q3	2.85	0	2000	0	4.0

and 1.4 MV/m at the bunch frequency, 625 Hz. Eight bunchers in the ring (see Fig.2), each 1 cm long, provide 205 longitudinal buckets to capture bunches that have  $\pm 1.2\%$  velocity spread.

#### Source and Intensity

Molecules from a source, typically at 310 m/s, must be decelerated for injection into a ring. A linear decelerator array for this purpose consists of a series of pulsed field regions. A molecule of CH<sub>3</sub>F will enter each region at zero field and decelerate upon exit at full field. This sequence is reversed for ND<sub>3</sub>. An a.g. pattern of static focusing lenses must be added to confine the beam transversely; for CH<sub>3</sub>F these must also overcome the vertical defocusing from fringe fields. 120 pulsed decelerator electrodes and an equal number of lenses are required making an array of 15.1 meter length. Bunches that emerge from this are 8 mm long spaced 48 mm. They are injected into 205 buckets formed by the bunchers in the ring. The injection kicker electrode is a section of arc guide fields that is turned off rapidly after the buckets in the ring are filled. In the deceleration process, emittances of the bunches in the three orthogonal momentum-displacement spaces are conserved. These become the acceptances for molecules at the source and are amply filled by a typical xenon-seeded source. A CH<sub>3</sub>F bunch will have  $1 \times 10^7$  molecules. At the 625 Hz bunch rate this gives a circulating current of  $6 \times 10^9$ molecules/sec.

#### DISCUSSION

Designing a storage ring lattice for a strong-seeking molecule is far more complex than that for a weak-seeking one. The arc sectors must have an alternating-gradient structure. The higher-order component in a lens field,  $A_5$  in this case, plays a crucial role in attaining a reasonable dynamic aperture. The vertical closed orbit due to gravity has to be corrected. However, in this initial effort, solutions are found.

### ACKNOWLEDGEMENTS

We thank Swapan Chattopadhyay and Ying Wu for early assistance with this work, and David Robin for useful discussions.

#### REFERENCES

- H. Bethlem, G. Berden, and G. Meijer, "Decellerating Neutral Dipolar Molecules" Phys. Rev. Lett. 83, 1558 (1999).
- [2] H. Nishimura, G. Lambertson, J. G. Kalnins, and H. Gould, "Feasibility of a synchrotron storage ring for neutral polar molecules" to be published in Rev. Sci. Instr. Preprint:LBNL-51597, http://arxiv.org/abs/physics/0212044
- [3] G. Lambertson, "Beam Dynamics in a Storage Ring for Neutral (Polar) Molecules" in this proceedings, Portland, May 2003.