THE MAGNETIC DESIGN AND PERFORMANCE OF THE SRRC-U9 UNDULATOR

S.C. Gottschalk, D.C. Quimby, K.E Robinson^{*} STI Optronics, Inc., Bellevue, WA 98004 USA

*Presently with Lawrence Berkeley National Laboratory.

Abstract

The magnetic design of the recently completed SRRC-U9 undulator system employs several enhancements to improve its performance. Side magnets to increase onaxis field augment the wedged-pole¹ magnet geometry. Partial volume and partial strength end magnets and tuning bars achieve a zero-displacement, rapid turn-on, end-field configuration.² Gap dependent normal and skew component ambient (earth) field compensation has been included in the design. The design includes bot active and passive end correction. Optical phase and trajectory shimming techniques have been employed to optimize performance. For k=10 the performance is greater than 80% of ideal through the seventh harmonic. At a K=6.4 the calculated spectral performance is greater than 80% through the thirteenth harmonic. A complete battery of magnetic measurements has fully established the field, spectral and beam dynamic characteristics of the undulator.

1 INTRODUCTION

This paper describes the magnetic design of the U9 undulator.³ It is a wedged pole hybrid planar insertion device built for the Synchrotron Radiation Research Company located in Hsinchu, Taiwan, R.O.C. A detailed discussion of the magnetic design method that we use has been described elsewhere.⁴ This device had three new features: 1) side magnets which were added to increase the peak field, 2) non-displacing entrance and exit built by using partial strength and reduced volume end magnets and 3) pre-tuning to remove ambient field differences between the STI Optronics site in Seattle, Washington, and SRRC.

2 CENTRAL FIELD

Wedged poles hybrids were used to generate the high fields required by SRRC with 20% less magnet volume than a straight pole. In addition, wedged magnets were small enough to permit single piece fabrication, unlike straight sided magnets. An isometric view of a quarter period cell is shown in Figure 1. The wedge angle, chamfers and magnet recess were chosen to remove saturated parts of the pole, reduce demagnetizing fields, allow sufficient space for tuning shims and control higher field harmonics.⁵ The design approach was conservative: low remanence magnets were used for peak field

estimates but high remanence was used for pole saturation, transverse rolloff and parts tolerancing. This had two ramifications: 1) side magnets are needed for field strength margin, and 2) pole transverse width was increased to 8.0cm to keep the transverse rolloff small (0.035% at ± 10 mm was calculated), in the event that stronger magnets were delivered by the vendor.

Adding side magnets reduced transverse flux leakage to 0.5% without degrading the transverse rolloff, changing field harmonics or reducing the pole permeability. For ease of fabrication we selected a side magnet that increased the field strength 5.2%. Larger side magnets could have increased the field by 8%. While the side magnets did have clamps, the lack of a transverse overhang caused them to be so strongly attracted to the poles that special tooling was needed to remove them during end field adjustment. Our simulated annealing magnet sorting algorithm was modified to include steering and multipole moments that can be created by side magnets. Only a small, easily removed quadrupole was observed on the untuned device. When side magnets are properly sorted, they do not appear to degrade field errors.

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| Quantity | Specification | MAGNET FEA | Experiment |
|----------------------------------|---------------|------------|------------|
| Peak field, g=18mm | 1.245 T | 1.365 T | 1.365 T |
| Transverse rolloff at 10mm | < 0.1% | 0.035% | 0.037% |
| 3 rd Harmonic content | <10% | 5.8% | 5.8% |



Figure 1: Isometric view of quarter period cell.

3 AMBIENT FIELD

Hybrid ID's can modify any ambient fields that are present causing gap dependent changes in beam dynamics which could require retuning the device at the new location. The following hypothetical example illustrates the point. Site A has a 0.2 Gauss ambient while site B has a 0.6 Gauss ambient. The ID 450cm long and has no 1st or 2^{nd} integral. Normal field enhancement is about 1.8 at minimum gap. For site A the ambient field inside the ID is increased to 1.8*0.2 = 0.36 Gauss and the 1st integral changes by (0.36-0.2)*450=72 G-cm while the 2nd has changed by 16,200 G-cm^2. The same ideal ID is moved to site B. Now the 1st integral changes by 216 G-cm as the gap is varied and the 2nd integral will appear to change by 48,600 G-cm². We analyzed these effects in advance and pre-tuned the device for the ambient field at SRRC. When we arrived at SRRC we found that the normal and skew integrals agreed to within 10 G-cm.

It is intuitively clear that at a small gap the ferromagnetic poles will shunt any skew ambient fields away from the e-beam and leave no on-axis skew field. Skew shunting will decrease as the gap is opened until it becomes absent for an infinite magnetic gap. Skew shunting is shown in Figure 2 where almost complete skew shunting can be seen at small gaps.

In the normal field direction, the field lines that would have gone through the magnets are shunted into the high permeability poles which leads to an overall normal field enhancement at small gaps. The field profiles for several gaps are shown in Figure 3. Skew and normal field effects after wiggle period averaging are summarized in Figure 4.



Figure 2. Skew field shunting at ends vs gap; gaps = 18 mm to 140 mm



Figure 3. Normal field enhancement at ends vs gap; gaps = 18 mm to 140 mm



Figure 4. Gap dependence of centerline shunting effects, normal and skew fields.

4 AIR GAP END CORRECTORS

The U9 device is equipped with normal field, fixed gap, air core end correctors shown in Figure 5. If SRRC needs to actively point the photon beam, use of the control software can energize the correctors in a gap dependent manner. The coils have 49 turns and a rectangular shape with the steering legs are placed at ± 30 degrees from vertical to eliminate sextupole moments. There is some loss in efficiency due to fringe fields being shunted into the first pole. The correctors did not introduce any significant multipole moments at their maximum currents of 10 amps. There are no skew field end correctors.



Figure 5. Air gap end corrector photo.

5 STRONGBACK SHAPING TO REDUCE GAP DEPENDENT PHASE ERRORS

For the U9 device there is a 12.5 metric ton variation in the attractive force between minimum and maximum gaps. Strongback deflections under magnetic loads will cause local gap variations which leads to a reduced spectral brightness. Spectral analyses for a collection of 25 previously built ID's having periods from 1.8cm to 8.5cm and K=11.9 to 0.5 were used to test the Walker phase error model.⁶ The model did an excellent job of predicting the *average* harmonic intensity loss caused by phase errors. This allowed us to predict phase errors based on local gap variations due to strongback deflections vs gap.

The ideal situation is to eliminate any deflection by using large, heavy strongbacks. For U9 there were overall height and weight restrictions. We used stainless steel strongbacks rather than aluminum and devised a new 4 point beam hanger system to reduce the deflection to 25 microns. This produced the gap dependent phase error variation of 2 degrees which would produce a 1.0% loss in 3^{rd} harmonic intensity.

Experimental phase error plots showed that strongback deflection under load was responsible for almost all of the non-ideal spectral performance. The remaining spectral error sources were several micron gap dependent tapers (<0.5 degrees) and field errors.

6. CONCLUSION

The SRRC/U9 met all magnetic requirements. It used side magnets to increase the on-axis field without any performance degradation. Ambient field correction succeeded in predicting changes to 1st and 2nd integrals and internal trajectories due to gap dependent difference in the ambient fields at STI and SRRC. The Walker phase error model⁶ was verified on an ensemble of 25 APS ID's having $\lambda_w = 18$ mm through 85 mm and was used to specify strongback deflection.

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