# **EXPERIMENTAL POLARIZATION CONTROL OF THOMSON SCATTER-ING X-RAY SOURCE**

Zhang Hongze<sup>†</sup> Tsinghua University, Beijing, China

## Abstract

maintain

must

## EXPERIMENT We carry out the X-ray polarization control and meas-

attribution to the author(s), title of the work, publisher, and DOI Thomson scattering of intense laser pulses from relativistic electrons allows us to generate high-brightness and tunable-polarization  $X/\gamma$ -ray pulses. In this paper, we demonstrate the polarization control of the Thomson scattering source experimentally. We control the incident laser polarization by rotating a quarter-wave plate, thus controlling X/ $\gamma$ -ray polarization. In order to measure the polarization, we use Compton scattering method. Meanwhile, stokes parameters of  $X/\gamma$ -ray whose energy varies between tens of keV and MeV are simulated. The simulation results show that with the increasing of X-ray Energy, X-ray polarization is a constant value in a small cone of motivation. According to modulation curves analysed from experiment results, we can get the conclusion that the polarization of Thomson scattering source is tunable and controllable.

### **INTRODUCTION**

distribution of this work Polarized X/ $\gamma$ -ray has been studied and widely used in various scientific field. Polarization of  $X/\gamma$ -ray can provide unique information addition to X-ray imaging and analysis of spectroscopy for researchers. In astrophysics, polarimetric observations of neutron stars provide the information of N the intensity and geometry of the magnetic field [1-4]. In material science and biology, polarized X-rays can enhance the sensitivity of X-ray fluorescence analysis [5]. In nu-201 clear physics, polarized  $\gamma$ -rays plays an important role in licence (© studying nuclear property. We can study the structure of nuclear by nuclear resonance fluorescence with polarized X/ $\gamma$ -ray (NFR) [6-9]. And the polarized state of  $\gamma$ -rays is 3.0 also important for the measurements of the parity of the nuclear states [10], the investigation of giant resonances of B nuclei and the scattering reactions between photons and the CC nuclei [11,12].

Compton scattering is the elastic scattering of a photon of from a free electron, for the low energy electron(~MeV), it is also called as Thomson scattering [13-15]. Compton(Thomson) scattering X-ray source has been studied under the and developed for decades [16-20]. Comparing to the mechanism of other radiation sources, it can produce ultrashort, energy continuously tunable, high brightness, wellcollimate and high polarized X-ray beams by laser photons scattering from free relativistic electrons [21-24]. Because é of the advantages in X-ray application, Thomson scattering mav X-ray source is utility in material, medical and biological work areas [24-28]. In our experiment, we change the polarization of X-ray by adjusting the polarization of laser beams from this since the polarization of laser is directly transferred to the scattered photons.

† zhz16@mails.tsinghua.edu.cn

urement experiment on Tsinghua Thomson scattering Xray source (TTX) platform. TTX is set up with a linac system and a femtosecond laser system. The linac system consists of a S-band photocathode RF gun, a magnet compressor and two x-band harmonic structures to generate high brightness electron pulse. The laser system can generate 266-nm ultraviolet pulse for the photocathode and 800-nm infrared pulse for the scattering interaction. The energy of X-ray photons is 50-keV and the flux is about  $10^7 s^{-1}$ [29]. In our experiment, laser photons track through the quarterwave plate and have a head-on interaction with the highquality electron beams in the vacuum interaction room. The polarization of scatted photons is determined by the incident laser beams, which are controlled by the quarterwave plate precisely.

In our experiment, we use the Compton scattering method, a kind of polarization-sensitive process which is more accurate than before, to measure the polarization of X-ray beams. According to the Klein-Nishina formula [30], for linear incident  $X/\gamma$ -ray photon, the azimuthal distribution of the scattered photons is strongly depended on the X-ray polarization. A target, made from polyethylene, is placed after the titanium window of the beam pipe. The size of the cylinder target is 5-cm in height and 0.75-cm in radius. X-ray pulses irradiate on the end of the cylinder and generate scattered photons. We use an image plate wrapping around the cylinder to record scattered photons (Figur 1). Meanwhile, we use two thin aluminium rings locked to both ends of the polyethylene cylinder to support the image plate. The curved image plate is 2.5-cm in radius.



Figure 1: Schematic and real picture of target.

## RESULTS

Figures 2-4 show the experiment results recorded by image plate and the simulation results done with Geant4.



Figure 2: Linear polarized incident laser beams. First subpicture: Simulation result of the scattered photons recorded by the image plate. Second subpicture: experiment result of the scattered photons recorded by the image plate. Third subpicture: Calculated results from the simulation (red) and experiment (green) results.



Figure 3: Ellipse polarize incident laser beams.



Figure 4: Circle polarized incident laser beams.

In Figs 2-4, for each polarization result, we sum the photons recorded by the image plate in the z-direction and get modulation curves. According to the modulation curves, we calculate the modulation amplitude and polarization of different polarized X-ray pulses. The modulation curves' amplitudes vary from the maximum value for linear polarization to the minimum value for circle polarization. The reasons, affecting the accuracy of the results, include the jitter of the electron pulse and laser pulse. Also, we assume that the X-ray bunches are parallel after tracking through the filter. But there is still a small divergence for each Xray bunch in the forward direction.

SAP2017, Jishou, China JACoW Publishing doi:10.18429/JACoW-SAP2017-WEBH3



Figure 5: Red line: the theory result of modulation amplitude. Green line: the simulation result of modulation amplitude. Red dot line: the simulation result of polarization. Red spots: experiment results of polarization. Blue spots: experiment results of modulation amplitude.

Simulated result (green line) is smaller than the theory result (red line). Because when we calculate the theory result, we only consider primary scattering process and ignore the background. The variation trend of the experiment result fits the simulation results. However, there is still one point doesn't fit the simulation results well. Because the simulation process is done under an ideal condition. The polarization of X-ray irradiated on the target is supposed to be the same. However, X-rays generated by Thomson scattering sources have a small divergence in forward direction. It results in the polarization of X-ray irradiated on the target varies between a small region. We can reduce the target's cross section radius to improve the accuracy of the experiment.

The X-ray energy is about 50-keV in our experiment and the X-ray polarization is nearly a constant value. In order to verify the polarization relationship between incident laser pulses, electron beams in high energy section, we calculate stokes parameters of X-rays under different electron beams energy. Limited by the experiment condition, we simulate the interaction process with Cain program. Stokes parameter  $S_2$  represents the circle polarization and  $S_3$  represents the linear polarization.  $S_2 = 1$  or  $S_2 = -1$  represents clockwise and counterclockwise. And  $S_1(S_3) = 1$  or  $S_1(S_3) = -1$  represents two polarized directions which are orthogonal to each other. Complete polarized states have  $S^2 = 1(S^2 = S_1^2 + S_2^2 + S_3^2)$ , but mixed states have  $S^2 < S_1^2 + S_2^2 + S_3^2$ 1. Changing energies of electron beams in Cain, the average value of  $S_2$  and  $S_3$  in a small cone ( $\theta = 1/5\gamma$ ) under different X-ray energys are shown in Fig. 6.



Figure 6: Stokes parameters. Green line: average value of S\_2 under different X-ray energy. Red line: average value of S\_3 under different X-ray energy.

DOI.

and I

attribution to the

tain

maint

must

work

Any distribution of this

BY 3.0 licence (© 2017).

In Fig. 5. Stokes parameters are almost invariant with Xray energy increasing in a constant small  $cone(\theta=1/5\gamma)$ . publisher, This means energies of electron beams doesn't affect the polarization of X-ray with a wide range of energy in a small cone. In order to measure the high-energy X-ray polarizawork, tion, a smaller cross section target is useful. But with high energy X-ray, about 10-Mey or more, the effect of pair prohe duction will affect the accuracy of the measurement. author(s), title of

#### CONCLUSION

Thomson Scattering Source is an important way to generate polarized  $X/\gamma$  -ray pulses which are tunable and controllable. Polarized X/ $\gamma$ -ray pulses are important and useful in various scientific area. Our experiment is the first time that try to precisely measure the polarization of Thomson Scattering Sources and verify the relationship of polarization between the incident laser beams and scattered X-ray beams. We use Compton scattering method, applicable in the energy range keV  $\sim$  MeV, to measure the polarization of X-ray beams. The experiment results show that we can produce accuracy polarization  $\gamma$ -ray/X-ray pulses by changing the polarized state of incident laser beams. And X-ray polarization are nearly constant value in a small cone under different X-ray energy.

#### REFERENCES

- [1] P. A. Connors and R. F. Stark, Nature 269, 8 (1977).
- [2] E. Costa, P. Soffitta, R. Bellazzini, A. Brez, N. Lumb and G. Spandre, Nature 411, 662 (2001).
- [3] J. R. P. Angel, R. Novick, P. V. Bout and R. Wolff, Phys. Rev. Lett. 22, 861(1969).
- [4] G. Bao, P. J. Witta and P. Hadrava, Phys. Rev. Lett. 77, 12(1996).
- [5] J. O. Christoffersson and S. Mattsson, Phys. Med. Biol. 28, 1135(1983).
- [6] L. I. Schiff, Phys Rev., 70, 761(1946).
- [7] U. Kneissl, H. H. Pitz, A. Zilges and Prog. Part. Nucl. Phys., 37, 349(1996).
- [8] M. E. Rose and R. L. Carovillano, Phys. Rev., 122, 1185 (1961).
- [9] C. T. Angell, R. Hajima, T. Hayakawa, T. Hayakawa, H. J. Karwowski, and J. Silano, Phys. Rev. C, vol. 90, 054315(2014).
- [10] N. Pietralla, Z. Berant et al., Phys. Rev. Lett., vol. 88, 012502 (2001).
- under the terms of the CC [11] H. Arenhövel, and E. Hayward, Phys. Rev., vol. 165, 1170(1968).
- used [12] E. Hayward, W. C. Barber and J. Sazama, Phys. Rev. C, ę vol. 8, 1065 (1973).
- nay <sup>2</sup>[13] A. H. Compton, *Phys. Rev.*, vol. 21, 483(1923).
- work [14] R. H. Milburn, Phys. Rev. Lett., vol. 10, 75(1963).
- [15] G. Fiocco, and E. Thompson, Phys. Rev. Lett., vol. 10, 89(1963).
- Content from this [16] F. R. Arutyunian, and V. A. Tumanian, Phys. Lett., vol. 4, 176 (1963).

- [17] A. M. Sandorfi et al., IEEE Tran. Nucl. Sci. vol. 30, 3083 (1983).
- [18] R. W. Schenlein et al., Science, vol. 274, 236(1996).
- [19] C. Bula et al, Phys. Rev. Lett., vol. 76, 3116(1996).
- [20] A. Ting et al., J. Appl. Phys, vol. 78(1), 575(1995).
- [21] V. N. Litvinenko et al., Phys. Rev. Lett., vol. 78, 4569 (1997).
- [22] I. V. Pogorelsky et al., Phys. Rev. ST Accel. Beams, vol. 3, 090702 (2000).
- [23] D. J. Gibson et al., Phys. Rev. ST Accel. Beams, vol. 13, 070703 (2010).
- [24] P. Oliva et al., Appl. Phys. Lett., vol. 97, 134104(2010).
- [25] K. Yamada et al., Nucl. Instr. Meth. A, vol. 608, S7 (2009).
- [26] R. Kuroda et al., Nucl. Instr. Meth. A, vol. 608, S28 (2009).
- [27] J. Abendroth et al., Journal of structural and functional genomics, vol. 11, 91 (2010).
- [28] M. Bech et al., Journal of synchrotron radiation 16, 43 (2009).
- [29] Y. Du et al., Nucl. Instr. Meth. A, vol. 637, S168 (2011).
- [30] L. W. Fagg, S. S. Hanna, Rev. Mod. Phys. vol. 31, 711(1959)

WEBH3